

# Transduction Mechanisms and Power Density for MEMS Inertial Energy Scavengers

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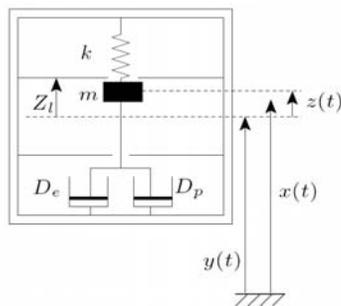
## Abstract

In this paper, the trends from the last 10 years of inertial micro-generator literature are investigated and it is shown that, although current generator designs are still operating well below their maximum power, there has been significant improvement with time. Whilst no clear conclusions could be drawn from reported fabricated devices with respect to preferred transducer technology, this paper presents operating charts for inertial micro-generators which identify optimal operating modes for different frequencies and normalized generator sizes, and allows comparison of the different transduction mechanisms as these parameters vary. It is shown that piezoelectric generators have a wider operating range at low frequency than electromagnetic generators, but as generator dimensions increase, the frequency to which piezoelectric transducers outperform electromagnetic transducers decreases.

*Keywords: Energy-harvesting, piezo-electric, ultra low-power*

## 1 - INTRODUCTION

Motion and vibration are attractive sources for micro-engineered energy scavenging generators [1, 2]. The most universal motion scavengers are of the inertial type, *i.e.* having a proof mass suspended within a frame, and energy extracted by a transducer that damps the motion of the proof mass within the frame. These devices have the advantage that they can function simply by being attached to a source of motion at a single point, rather than relying on the relative motion of different parts of the “host” structure. Thus they are also well suited to miniaturisation.



**Figure 1** - Schematic construction of inertial generators.

The basic operating principle of inertial micro-generators is illustrated in Fig. 1. The fundamental parameters limiting the generator output are its proof mass  $m$  and maximum internal displacement  $Z_l$ , and the source motion amplitude  $Y_0$  and frequency  $\omega$  [3]. From these we can derive the maximum power from basic principles. If we assume harmonic source motion, the maximum acceleration  $a_{max}$  is  $\omega^2 Y_0$ . The

maximum damping force by which energy can be extracted is equal to the inertial force on the proof mass,  $ma_{max}$  (if greater, the mass will not move). If energy is extracted in both directions, and the internal motion amplitude  $Z_o = Z_l$ , (giving the maximum travel range of  $2Z_l$ ) we derive a total energy per cycle of  $4Z_l ma_{max} = 4Z_l m \omega^2 Y_0$ . To convert this to power we simply divide by the excitation period  $2\pi/\omega$ , giving:

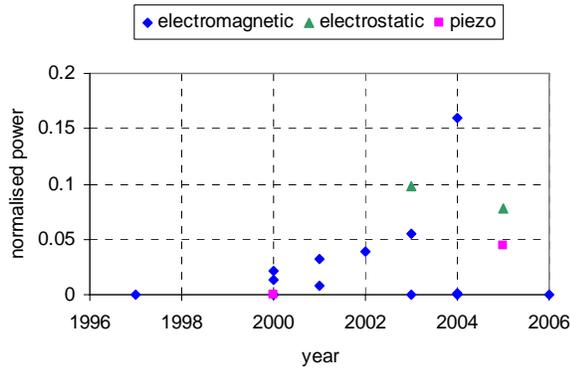
$$P_{max} = 2Y_0 Z_l \omega^3 m / \pi \quad (1)$$

We can then define a normalised power  $P_n = P/P_{max}$  as a measure of how close the performance of a specific device comes to the optimum level. We have calculated  $P_n$  for *measurements* on inertial energy scavengers reported in the literature [1, 4-25] and the resulting values are plotted in Fig. 2 as a function of year of publication. An upwards trend can clearly be seen, although the best values are still below 20% of the optimum. Although  $P_n$  should not drop with volume, since it is normalised to device size, the same data plotted against device volume (Fig. 3) show that typically the best  $P_n$  values have been achieved for larger devices. This is likely an indication of the technological difficulties encountered at smaller size scales, for example the greater difficulty in achieving high magnetic flux gradients. Finally, we plot the normalised power vs. frequency (Fig. 4), and a downwards trend is clearly seen.

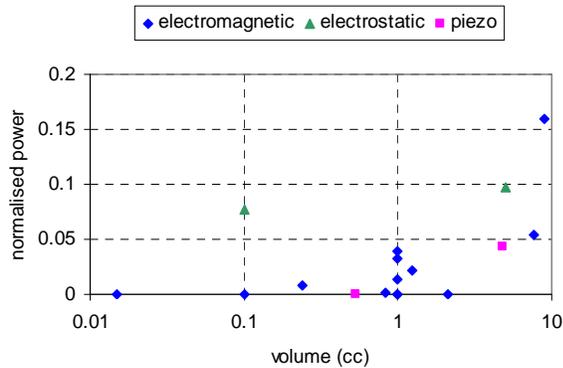
Several transduction mechanisms can be used for inertial micro-generators, namely electromagnetic [14], electrostatic [26] and piezoelectric [27]. The transduction type is also indicated in Figs. 2 – 4, but no clear trends can be seen regarding their relative merits. To obtain some general guidance on the practical limitations of specific transduction methods, we have examined the key issue of the damping levels that can be achieved. We consider only mechanically

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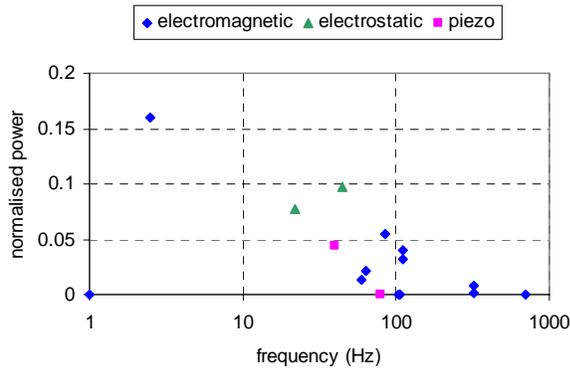
resonant devices operating at the resonant frequency  $\omega_n$ , as this covers most reported examples.



**Figure 2** - Normalised measured power  $P_n$  vs. year of publication.



**Figure 3** - Normalized measured power  $P_n$  vs. device volume.

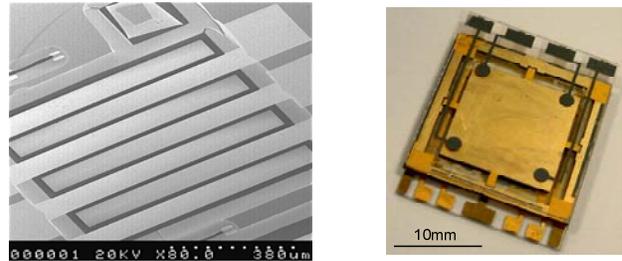


**Figure 4** - Normalised measured power  $P_n$  vs. operating frequency.

## 2 - SCALING ANALYSIS OF TRANSDUCER TYPES

In an ideal case, the parasitic damping would be zero, and maximum power is then obtained by setting the electrical damping factor  $\zeta_e$  to the level that allows the mass to move over the entire internal range, but without hitting the end-stops, i.e.  $Z_0=Z_l$  [3]. However, in some cases the maximum damping force of the transducer is insufficient to achieve this, and thus the device cannot operate in a resonant mode. In cases where the parasitic damping factor  $\zeta_p$  is not negligible,

maximum power will inevitably be reduced, but the optimum  $\zeta_e$  will still be that which gives  $Z_0=Z_l$ , unless this requires  $\zeta_e < \zeta_p$ , in which case  $\zeta_e = \zeta_p$  should be chosen if possible.



**Figure 5** - Prototype MEMS inertial scavengers: piezoelectric (left) and electrostatic (right).

In general, the damping factor is related to the resonant quality factor  $Q$  by  $Q=1/2\zeta$ . Since the damping will have parasitic and electrical (transducer) components, we introduce the quantities  $Q_p = 1/2\zeta_p$  and  $Q_e = 1/2\zeta_e$ . Note that the combined  $Q$  is given by  $1/Q = 1/Q_p + 1/Q_e$ . Furthermore, for a resonant system in which the damping force is proportional to the relative mass-frame velocity, i.e.  $F=-Ddz/dt$ ,  $D = 2m\omega_n\zeta$ . To perfectly damp the system we require  $Q = Z_l/Y_0$ , so that  $D = m\omega_n/(Z_l/Y_0)$ .

An electromagnetic damper can be implemented as a coil moving across the boundary of a region of magnetic flux density  $B$ . If we assume that the induced voltage is limited mainly by the resistive load  $R$ , rather than the coil's inductance, then in this case the electrical damping coefficient is given by  $D_e = (NBl)^2/R$  [3], where  $N$  is the number of coil turns and  $l$  is the length of the border of the flux region cutting across the coil. The load  $R$  will consist of the coil resistance  $R_i$  (which is a parasitic component) and the energy extracting load  $R_L$  in series. The former determines the maximum  $D_e$ , and thus the minimum  $Q_e$ ; however, unless  $R_L$  is substantially greater than  $R_i$  most of the electrical power will be wasted. Instead we assume that a useful device has at least  $R = 10R_i$ . This sets the minimum achievable  $Q_e$  as:

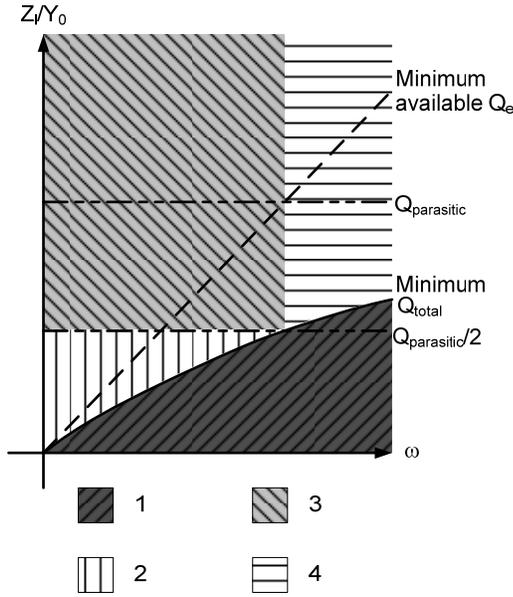
$$Q_e(\min) = \frac{10mR_i}{(NBl)^2} \omega_n \quad (2)$$

If this quantity is  $> Z_l/Y_0$ , it will not be possible to achieve harmonic motion by electrical damping alone, although it may be if significant parasitic mechanical damping is present.

A similar quantity can be derived for the piezoelectric case. For a piezo element of capacitance  $C_0$  the maximum electrical damping coefficient is obtained for an electrical load  $R_L=1/\omega C_0$ , and is given by  $D_e = \alpha^2/2\omega C_0$  [28]. Here  $\alpha = e_{33}A/t$ , with  $A$  and  $t$  the cross-sectional area and thickness of the piezo element, and  $e_{33}$  the piezoelectric coefficient. In practice, since piezo elements can only achieve very small direct displacements, devices must incorporate some leverage mechanisms (of ratio  $r$ ). This will transform the damping coefficient by a factor of  $r^2$ . Thus we obtain:

$$Q_e(\min) = \frac{2mC_o r^2}{\alpha^2} \omega_n^2 \quad (3)$$

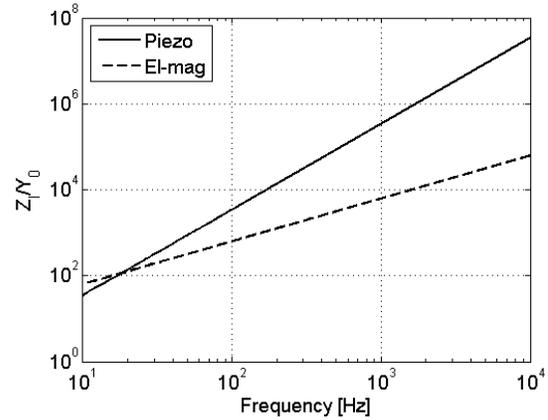
If we define  $Z_l/Y_0$  and  $\omega_n$  as our general operating parameters, then we can plot (2) and (3) as limits of operation for electromagnetic and piezoelectric devices. By adding the  $Q_p$  of the device, and the combined  $Q$ , we can indicate regions of operation as a function of operating point.



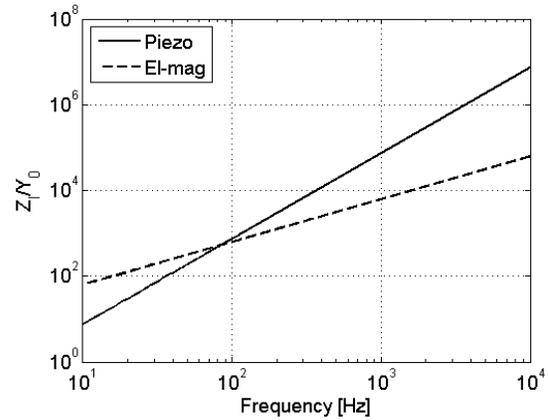
**Figure 6** - Operating limits on electrical damping and parasitic Q factor.

This is done in Fig. 6 schematically, for an electromagnetic generator with parasitic damping present. The parasitic damping factor (and hence the  $Q_p$ ) has been assumed to be constant with frequency for simplicity, but any actual frequency dependence could easily be substituted. The operating regions are as follows:

1. Harmonic motion is not possible. The maximum combined damping is less than required for oscillation within  $Z_l$  and so the mass will strike the end-stops.
2. For optimal operation, the electrical damping should be set so that the  $Z_0=Z_l$ . Electrical damping will be greater than parasitic damping, and so this generator can achieve the maximum power for the level of parasitic damping present.
3. The electrical damping should be set to equal the parasitic mechanical damping, i.e.  $Q_e=Q_p$ . This will give  $Z_0/Y_0 = Q_{\text{parasitic}}/2$ .
4. The electrical damping should be set to the maximum that can be achieved, but it will still be less than the parasitic mechanical damping, i.e.  $Q_e > Q_p$ . The generator can operate within the displacement constraint, but a different transducer could in principle extract more power.



**Figure 7** - Comparison of minimum Q factors with electromagnetic and piezoelectric cube devices of volume 1cc



**Figure 8** - Comparison of minimum Q factors with electromagnetic and piezoelectric cube devices of volume 0.1cc.

Figures 7 and 8 show two specific examples of the minimum Q-factor achievable from electromagnetic and piezoelectric generators. In each case we assume a cubic device of length  $L$  and the mass, of relative density 8.9 (Ni) is taken to occupy half the device volume. The electromagnetic device is assumed to have a flux density of 1 T and a copper coil occupying 2% of the device volume, for which a fixed  $N^2/R_i$  is obtained. The active coil length  $l$  is assumed to be  $L/2$ . For the piezoelectric device, we assume  $\epsilon_r = 1000$ , area  $L^2$ , thickness  $L/10$  and  $e_{33} = 0.15 \text{ C/m}^2$ . A leverage factor of 500 was chosen. Because the  $Q_{\min}$  for electromagnetic and piezoelectric devices scale as  $\omega$  and  $\omega^2$  respectively, there will always be a frequency above which electromagnetic devices can achieve higher damping. As can be seen in Figure 7 and Figure 8, the cross-over frequency increases as device size decreases. It can also be observed that the increasing Q with frequency (for both transducer types) could explain the decreasing performance trend seen in Fig. 4.

### 3 - CONCLUSIONS

Obtaining maximum power from inertial energy scavengers is often limited by the maximum damping force achievable in the transduction mechanism, and this problem increases with increasing frequency. Piezoelectric generators can outperform electromagnetic generators at low frequency, but with increasing frequency, the internal capacitance of the piezoelectric reduces the amount of real power that can be obtained. This suggests that piezoelectric devices might be better suited to human body powered applications and electromagnetic devices to high frequency applications.

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